薄壁构件试验模型的动力学相似设计方法

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摘要:针对由薄壁构件相似模型试验结果预测原型动力学特性的问题,建立了统一的数学模型,并基于方程分析法推导了动 力学相似关系。通过对薄壳单元进行受力分析,推导得到薄壁构件的微分方程,并通过微分方程得到完全几何相似模型可准确预测 原型动力学特性的相似关系。在建立薄壁构件结构参数对固有频率的敏感性与相似因子指数的比值关系基础上,提出薄壁构件不 完全几何相似模型与原型的动力学相似关系(即畸变相似关系)的确定方法。最后给出基于敏感性分析的薄壁构件精确畸变相似关 系设计流程,为薄壁构件相似试验模型的设计及动力学特性的预测提供参考。

关键词:薄壁构件;相似设计;畸变相似关系;敏感性;试验模型

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Dynamic Similitude Design Method of Experimental Models for Thin Walled Structures

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Abstract: Aiming at the problem that experimental results of elastic thin walled structures similitude models predict dynamic characteristics of the prototype, the unified mathematical model was established and dynamic scaling laws was proposed based on the equation analysis method. The unified governing equation of elastic thin walled structures was firstly obtained by the stress analysis of a shell element, and the scaling law was deduced between the geometrically complete model and the prototype based on the unified governing equation. On the basis of establishing the ratio relationship between indexes of scaling factors and sensitivity results of structural parameters for the natural frequency, and the method of determining the dynamic scaling law (distorted scaling law) was proposed between the geometrically partial similitude models and the prototype. Finally, the design procedure of determining accurate distorted scaling laws was given out based on the sensitivity analysis, which provided the design method of similitude models and the prediction of dynamic characteristics for elastic thin walled structures.

Key words: thin walled structures; similitude design; distorted scaling law; sensitivity; test model

0 引言

薄壁构件通常是指厚度与结构件最小平面跨度 之比在 1/80 和 1/5 之间的弹性结构件^{III}。在工程应用 中,由于薄壁结构件具有结构简单、抗弯刚度大等优 点,被广泛应用于航空航天、海洋机械、化工机械等工 程领域^{III}。在工作过程中,薄壁构件受力情况复杂且经 常处于振动状态,其剧烈振动可能导致结构损坏,甚 至引起无法预料的破坏,因此,对其进行振动特性分 析具有重要意义^[3]。然而在实际研究中,若直接采用原 型进行试验就会受到体积大、试验难度大、成本高等 多方面因素的限制。因此,设计相似模型预测原型的 动力学特性具有重要价值。有很多学者对薄壁构件的 振动特性进行了研究;Leissa^[4]对薄壳的振动特性进行

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了理论分析,并给出了薄壁壳的振动微分方程: Narita[®]通过采用改进的 Ritz 法研究了简支板的固有 特性; Irie 等¹⁰分析各类层合壳的动力学特性,并提出 了用传递矩阵法分析薄壳在任意边界下的振动特性; Zhou 等ⁿ根据混合能量守恒原理,采用 Hamiltonian 方 法分析了薄圆板和圆环薄板的固有特性,得到了不同 边界下的频率方程;对于薄壁构件的相似设计方面的 研究,Krayterman 等¹⁸采用量纲分析法推导了薄板的 相似关系,考虑了边界条件对相似关系的影响,并对 相似模型设计方法的预测精度进行了讨论;Qian 等¹⁹ 得到了响应下层合板的相似关系,并且相似准则能准 确描述原型的响应:Rezaeepazhand 和 Simitses¹⁰通过 相似模型预测了层合壳的屈曲和自由振动特性; Ungbhakorn 和 Singhatanadgid^[11-12]提出了 1 种推导相 似关系的新方法,并且呈现了对称层合板和层合圆柱 壳在考虑载荷影响下的相似关系; Oshiro 和 Alves^[13]研 究了预测原型动力学特性的畸变相似关系,并且分析 了受动载荷原型的3个问题。

综上所述,对于典型薄壁构件,如弹性薄板、薄壁 圆柱壳的相似性研究也有很多,然而确定薄壁构件精 确的畸变相似关系研究尚不多见。

本文通过微元法给出了薄壁构件统一形式的本 构方程,并提出了确定薄壁构件精确畸变相似关系的 方法,为航空发动机等重大机械装备薄壁构件相似试 验模型的设计提供了参考。

1 基本理论

在笛卡尔坐标系 Oxyz¹ 中,建立曲线坐标系 O' $\alpha\beta\gamma$ 如图 1 所示。 α 、 β 沿曲面的 曲率线方向, γ 与 α 、 β 垂 直,位移 u、v、w 分别代表 α 、 β 、 γ 方向的切向位移。dr_o 为曲面上的微段弧长。E 为 材料的弹性模量, μ 为泊 松比, ρ 为密度。



图 1 曲线坐标系

曲面上的微段弧长 dr 可表示为^[14] dr= $A^2(d\alpha)^2+B^2(d\beta)^2$ (1) 式中:A、B分别为曲面拉梅参数,

$$\mathbf{A} = \sqrt{\left(\frac{\partial \mathbf{X}}{\partial \alpha}\right)^2 + \left(\frac{\partial \mathbf{Y}}{\partial \alpha}\right)^2 + \left(\frac{\partial \mathbf{Z}}{\partial \alpha}\right)^2}$$

$$B = \sqrt{\left(\frac{\partial \mathbf{x}}{\partial \beta}\right)^{2} + \left(\frac{\partial \mathbf{y}}{\partial \beta}\right)^{2} + \left(\frac{\partial \mathbf{z}}{\partial \beta}\right)^{2}} \circ$$

薄壁构件的内力与中面位移关系式为
$$\begin{cases} \mathsf{N}_{\alpha} = \mathsf{K}(\varepsilon_{\alpha}^{0} + \mu \varepsilon_{\beta}^{0}) \\ \mathsf{N}_{\beta} = \mathsf{K}(\varepsilon_{\beta}^{0} + \mu \varepsilon_{\alpha}^{0}) \\ N_{\alpha\beta} = N_{\beta\alpha} = \mathsf{K}\frac{(1-\mu)}{2}\varepsilon_{\alpha\beta}^{0} \\ \mathsf{M}_{\alpha} = \mathsf{D}(\chi_{\alpha} + \mu \chi_{\beta}) \\ \mathsf{M}_{\beta} = \mathsf{D}(\chi_{\beta} + \mu \chi_{\alpha}) \\ \mathsf{M}_{\alpha\beta} = \mathsf{M}_{\beta\alpha} = \mathsf{D}\frac{(1-\mu)}{2}\chi_{\alpha} \end{cases}$$

$$(2)$$

式中:K= $\frac{Eh}{1-\mu^2}$ 为薄膜刚度;D= $\frac{Eh^3}{12(1-\mu^2)}$ 为弯曲刚度,; $\varepsilon_{\alpha}^{0}, \varepsilon_{\beta}^{0}, \varepsilon_{s}^{0}$ 为中曲面的薄膜应变分量; $\chi_{\alpha}, \chi_{\beta}, \chi_{\alpha\beta}$ 为中曲面弯曲应变—曲率变化分量,可表示为

$$\begin{cases} \varepsilon_{\alpha}^{0} = \frac{1}{A} \frac{\partial u}{\partial \alpha} + \frac{v}{AB} \frac{\partial A}{\partial \beta} + \frac{w}{R\alpha} \\ \varepsilon_{\beta}^{0} = \frac{1}{B} \frac{\partial u}{\partial \beta} + \frac{u}{AB} \frac{\partial B}{\partial \alpha} + \frac{w}{R_{\beta}} \\ \varepsilon_{\alpha\beta}^{0} = \frac{A}{B} \frac{\partial}{\partial \beta} \left(\frac{u}{A}\right) + \frac{B}{A} \frac{\partial}{\partial \alpha} \left(\frac{u}{B}\right) \\ \chi_{\alpha} = \frac{1}{A} \frac{\partial}{\partial \alpha} \left(\frac{u}{R_{\alpha}} - \frac{1}{A} \frac{\partial \omega}{\partial \alpha}\right) + \frac{1}{AB} \left(\frac{u}{R_{\beta}} - \frac{1}{B} \frac{\partial \omega}{\partial \beta}\right) \frac{\partial A}{\partial \beta} \\ \chi_{\beta} = \frac{1}{B} \frac{\partial}{\partial \beta} \left(\frac{u}{R_{\beta}} - \frac{1}{A} \frac{\partial \omega}{\partial \beta}\right) + \frac{1}{AB} \left(\frac{u}{R_{\alpha}} - \frac{1}{B} \frac{\partial \omega}{\partial \alpha}\right) \frac{\partial A}{\partial \alpha} \\ \chi_{\alpha\beta} = \frac{A}{B} \frac{\partial}{\partial \beta} \left(\frac{u}{AR_{\alpha}} - \frac{1}{A^{2}} \frac{\partial \omega}{\partial \alpha}\right) + \frac{B}{A} \frac{\partial}{\partial \alpha} \left(\frac{v}{BR_{\beta}} - \frac{1}{B^{2}} \frac{\partial \omega}{\partial \beta}\right) \\ \vec{x} \oplus : R_{\alpha} \pi R_{\beta} \mathcal{H} \pm \vec{m} \approx \mathcal{H} \equiv 0 \end{cases}$$
(3)

将所有力矩向量分别在 $\alpha_{\lambda}\beta_{\lambda}\gamma$ 方向投影及所有 力向量分别对 $\alpha_{\lambda}\beta_{\lambda}\gamma$ 轴取矩,可得

$$\frac{\partial (BN_{\alpha})}{\partial \alpha} + \frac{\partial (AN_{\beta\alpha})}{\partial \beta} - \frac{\partial B}{\partial \alpha} N_{\beta} + \frac{\partial A}{\partial \beta} N_{\alpha\beta} + AB \frac{Q_{\alpha}}{R_{\alpha}} + ABq_{\alpha} - AB\rho h \frac{\partial^{2}u}{\partial t^{2}} = 0 \frac{\partial (AN_{\beta})}{\partial \beta} + \frac{\partial (BN_{\alpha\beta})}{\partial \beta} - \frac{\partial B}{\partial \alpha} N_{\alpha} + \frac{\partial B}{\partial \alpha} N_{\beta\alpha} + AB \frac{Q_{\beta}}{R_{\beta}} + ABq_{\beta} - AB\rho h \frac{\partial^{2}v}{\partial t^{2}} = 0 \frac{\partial (BQ_{\alpha})}{\partial \alpha} + \frac{\partial (AQ_{\beta})}{\partial \beta} - AB \frac{N_{\alpha}}{R_{\alpha}} - AB \frac{N_{\beta}}{R_{\beta}} + ABq_{y} - AB\rho h \frac{\partial^{2}W}{\partial t^{2}} = 0$$
(4)

式中:Q_α、Q_β为剪切力,有如下关系

$$\begin{cases} \frac{1}{AB} \left(\frac{\partial (AM_{\beta})}{\partial \beta} + \frac{\partial (BM_{\alpha\beta})}{\partial \alpha} - \frac{\partial A}{\partial \beta} M_{\alpha} + \frac{\partial B}{\partial \alpha} M_{\beta\alpha} \right) = Q_{\beta} \\ \frac{1}{AB} \left(\frac{\partial (BM_{\alpha})}{\partial \alpha} + \frac{\partial (BM_{\beta\alpha})}{\partial \beta} - \frac{\partial B}{\partial \alpha} M_{\beta} + \frac{\partial A}{\partial \beta} M_{\alpha\beta} \right) = Q_{\alpha} (5) \\ N_{\alpha\beta} - N_{\beta\alpha} + \frac{M_{\alpha\beta}}{R_{\alpha}} - \frac{M_{\beta\alpha}}{R_{\beta}} = 0 \\ \text{K} \quad \text{If } M_{\alpha\beta} + \frac{\partial (AN_{\beta\alpha})}{\partial \alpha} + \frac{\partial (AN_{\beta\alpha})}{\partial \beta} - \frac{\partial B}{\partial \alpha} N_{\beta} + \frac{\partial A}{\partial \beta} N_{\alpha\beta} + \\ \frac{1}{R_{\alpha}} \left[\frac{\partial (BM_{\alpha})}{\partial \alpha} + \frac{\partial (AM_{\beta\alpha})}{\partial \beta} - \frac{\partial B}{\partial \alpha} M_{\beta} + \frac{\partial A}{\partial \beta} M_{\alpha\beta} \right] - (6a) \\ AB\rho h \frac{\partial^{2} u}{\partial t^{2}} = 0 \\ \frac{\partial (AN_{\beta})}{\partial \beta} + \frac{\partial (BN_{\alpha\beta})}{\partial \alpha} - \frac{\partial A}{\partial \beta} N_{\alpha} + \frac{\partial B}{\partial \alpha} M_{\beta\alpha} + \\ \frac{1}{R_{\beta}} \left[\frac{\partial (AM_{\beta})}{\partial \beta} + \frac{\partial (BM_{\alpha\beta})}{\partial \alpha} - \frac{\partial A}{\partial \beta} M_{\alpha} + \frac{\partial B}{\partial \alpha} M_{\beta\alpha} \right] \\ AB\rho h \frac{\partial^{2} v}{\partial t^{2}} = 0 \\ \frac{\partial (AN_{\beta})}{\partial \beta} + \frac{\partial (BM_{\alpha\beta})}{\partial \beta} - \frac{\partial A}{\partial \alpha} M_{\alpha} + \frac{\partial B}{\partial \alpha} M_{\beta\alpha} \right] + \\ \frac{\partial (AN_{\beta})}{\partial \beta} + \frac{\partial (BM_{\alpha\beta})}{\partial \beta} - \frac{\partial A}{\partial \alpha} M_{\alpha} + \frac{\partial B}{\partial \alpha} M_{\beta\alpha} \right] + \\ \frac{\partial (AN_{\beta})}{\partial \beta} + \frac{\partial (BM_{\alpha\beta})}{\partial \beta} - \frac{\partial A}{\partial \alpha} M_{\alpha} + \frac{\partial B}{\partial \alpha} M_{\beta\alpha} \right] + \\ \frac{\partial (AN_{\beta})}{\partial \beta} + \frac{\partial (BM_{\alpha\beta})}{\partial \alpha} - \frac{\partial A}{\partial \alpha} M_{\alpha} + \frac{\partial B}{\partial \alpha} M_{\beta\alpha} \right] + \\ \frac{\partial (AN_{\beta})}{\partial \beta} + \frac{\partial (BM_{\alpha\beta})}{\partial \beta} - \frac{\partial A}{\partial \alpha} M_{\alpha} + \frac{\partial B}{\partial \beta} M_{\beta\alpha} \right] - (6c) \\ AB \left(\frac{N_{\alpha}}}{R_{\alpha}} + \frac{N_{\beta}}}{R_{\beta}} \right) - AB\rho h \frac{\partial^{2} \omega}{\partial t^{2}} = 0 \end{cases}$$

对于典型薄壁构件,如矩形薄板、圆环薄板、薄壁 短圆柱壳等,拉梅参数为实常数。所以在式(6)中,与 拉梅参数A,B对α,β的导数的相关项可以忽略。

另外,将式(2)、(3)代入式(6),可得薄壁构件的 本构方程

$$-\frac{BD}{A^{2}R_{\alpha}}\frac{\partial^{2}W}{\partial\alpha^{3}} - \frac{D}{BR_{\alpha}}\frac{\partial^{3}W}{\partial\alpha\partial\beta^{2}} + \frac{B}{A}\left(K + \frac{D}{R_{\alpha}^{2}}\right)\frac{\partial^{2}u}{\partial\alpha^{2}} + \frac{A(1-\mu)}{2B}\left(K + \frac{D}{R_{\alpha}^{2}}\right)\frac{\partial^{2}u}{\partial\beta^{2}} + (1+\mu)\left(\frac{K}{2} + \frac{D}{2R_{\alpha}R_{\beta}}\right)\frac{\partial^{2}v}{\partial\alpha\partial\beta} + (7a)$$

$$BK\left(\frac{1}{R_{\alpha}} + \frac{\mu}{R_{\beta}}\right)\frac{\partial w}{\partial\alpha} - AB\rho h\frac{\partial^{2}u}{\partialt^{2}} = 0$$

$$-\frac{D}{AR_{\beta}}\frac{\partial^{3}w}{\partial\alpha^{2}\partial\beta} - \frac{AD}{B^{2}R_{\beta}}\frac{\partial^{3}w}{\partial\beta^{3}} + \frac{(1+\mu)}{2}\left(K + \frac{D}{R_{\alpha}R_{\beta}}\right)\frac{\partial^{2}u}{\partial\alpha\partial\beta} + \frac{B(1-\mu)}{2A}\left(K + \frac{D}{2R_{\beta}^{2}}\right)\frac{\partial^{2}v}{\partial\alpha^{2}} + \frac{A}{B}\left(K + \frac{D}{R_{\beta}^{2}}\right)\frac{\partial^{2}v}{\partial\beta^{2}} + (7b)$$

$$AK\left(\frac{\mu}{R_{\alpha}} + \frac{1}{R_{\beta}}\right)\frac{\partial w}{\partial\alpha^{2}\partial\beta^{2}} - AB\rho h\frac{\partial^{2}u}{\partialt^{2}} = 0$$

$$-\frac{BD}{A^{3}}\frac{\partial^{4}w}{\partial\alpha^{4}} - \frac{2\mu D}{AB}\frac{\partial^{4}w}{\partial\alpha^{2}\partial\beta^{2}} - \frac{2D(1-\mu)}{AB}\frac{\partial^{4}w}{\partial\alpha^{2}\partial\beta^{2}} - \frac{2D(1-\mu)}{A}\frac{\partial^{4}w}{\partial\alpha^{2}\partial\beta^{2}} - \frac{2D(1-\mu)}{A}\frac{\partial^{4}w}{\partial\alpha^{2}\partial\beta^{2}} - \frac{2D(1-\mu)}{A}\frac{\partial^{4}w}$$

$$\frac{AD}{B^{3}} \frac{\partial^{4}w}{\partial\beta^{4}} + \frac{BD}{A^{2}R_{\alpha}} \frac{\partial^{3}u}{\partial\alpha^{3}} + \frac{D}{BR_{\alpha}} \frac{\partial^{3}u}{\partial\alpha\partial\beta^{2}} \frac{D}{AR_{\beta}} \frac{\partial^{3}v}{\partial\alpha^{2}\partial\beta} \quad (7c)$$

$$\frac{AD}{B^{2}R_{\beta}} \frac{\partial^{3}v}{\partial\beta^{3}} \left(\frac{BK}{R_{\alpha}} + \frac{\mu BK}{R_{\beta}}\right) \frac{\partial u}{\partial\alpha} + \left(\frac{\mu AK}{R_{\alpha}} + \frac{AK}{R_{\beta}}\right) \frac{\partial v}{\partial\beta} + \left(\frac{ABK}{R_{\alpha}^{2}} + \frac{2\mu ABK}{R_{\alpha}R_{\beta}} + \frac{ABK}{R_{\beta}^{2}}\right) w - AB\rho h \frac{\partial^{2}w}{\partialt^{2}} = 0$$

$$\frac{B}{D} = \frac{1}{2} \frac{1}{2} \frac{D}{D} \frac{\partial v}{\partial t} + \frac{1}{2} \frac{D}{D} \frac{D}{D} \frac{\partial v}{\partial t} + \frac{1}{2} \frac{D}{D} \frac{D}{D} \frac{\partial v}{\partial t} + \frac{1}{2} \frac{D}{D} \frac{\partial v}{\partial t} + \frac{1}{2} \frac{D}{D} \frac{\partial v}{\partial t} + \frac{1}{2} \frac{D}{D} \frac{D}{D} \frac{\partial v}{\partial t} + \frac{1}{2} \frac{D}{D} \frac{D}{D} \frac{D}{D} \frac{\partial v}{\partial t} + \frac{1}{2} \frac{D}{D} \frac{D}$$

由上式可知, 位移 u、v、w 对 α, β 的最高阶导数为 4。因此, 适用于薄壁构件统一的微分方程可归纳为

$$\sum_{j=u,v,w}\sum_{k=1}^{m=4}\sum_{i=0}^{n=k}L_{j-ki}\frac{\partial^{k}j}{\partial\alpha^{k-i}\partial\beta^{i}}+\sum_{j=u,v,w}L_{j}\frac{\partial^{2}j}{\partial t^{2}}=0 \qquad (8)$$

式中,j代表位移u、v、w;L_{j+ki}、L_j为相应的系数;t为时间。

- 2 相似关系
- 2.1 完全几何相似关系 原型和模型的本构方程为

$$\sum_{j=u,v,w} \sum_{k=1}^{m=4} \sum_{i=0}^{m=k} L_{jp-ki} \frac{\partial^k j_p}{\partial \alpha^{k-i} \partial \beta^i} + \sum_{j=u,v,w} L_{jp} \frac{\partial^2 j_p}{\partial t^2} \qquad (9a)$$

$$\sum_{j=u,v,w} \sum_{k=1}^{m=4} \sum_{i=0}^{n=k} L_{jm-ki} \frac{\partial^{k} j_{m}}{\partial \alpha^{k-i} \partial \beta^{i}} + \sum_{j=u,v,w} L_{jm} \frac{\partial^{2} j_{m}}{\partial t^{2}}$$
(9b)

式中:下标 p 代表原型, m 代表模型。

位移方程可表示为

$$\mathbf{j}(\alpha, \boldsymbol{\beta}, \mathbf{t}) = \mathbf{J}(\alpha, \boldsymbol{\beta}) \mathbf{e}^{\mathbf{i}\omega \mathbf{t}}$$
(10)

式中: ω 是固有频率;J表示模态函数 U、V、W。

将相似因子
$$\lambda_{e} = \frac{e_{p}}{e_{m}} (e=j, L_{j-ki}, \alpha, \beta, L_{j}, J) 代入原型的$$

本构方程(9a),可得

$$\sum_{j=u,v,w} \sum_{k=1}^{4} \sum_{i=0}^{k} \lambda_{Lj-ki} \frac{\lambda_{j}}{\lambda_{\alpha}^{k-i} \lambda_{\beta}^{i}} L_{jm-ki} \frac{\partial^{k} J_{m}}{\partial \alpha_{m}^{k-i} \partial \beta_{m}^{i}} = \sum_{j=u,v,w} \lambda_{Lj} \lambda_{j} \lambda_{\omega}^{2} L_{jm} J_{m} e^{i\omega t}$$
(11)

根据相似理论,原型与模型的本构方程中的系数 对应成比例^{16]},即

$$\lambda_{l_{j}-k_{i}} \frac{\lambda_{J}}{\lambda_{\alpha}^{k-i} \lambda_{\beta}} = \lambda_{l_{j}} \lambda_{J} \lambda_{\omega}^{2}$$
(12)

在完全几何相似的条件下

$$\lambda_{\alpha} = \lambda_{\beta} = \lambda_{h} = \lambda, \frac{\lambda_{L_{j,ki}}}{\lambda_{L_{j}}} = \frac{\lambda_{D}}{\lambda_{\rho}\lambda_{h}} \frac{1}{\lambda^{4}} = \frac{\lambda_{K}}{\lambda_{\rho}\lambda_{h}} \frac{1}{\lambda^{2}} = \frac{\lambda_{E}}{\lambda_{\rho}} \frac{1}{\lambda^{2}} (13)$$

因此, 薄壁构件的完全几何相似关系为

$$\lambda_{\omega} = \frac{1}{\lambda} \sqrt{\frac{\lambda_{\rm E}}{\lambda_{\rho}}} \tag{14}$$

2.2 畸变相似关系

通常情况下,直接采用完全几何相似模型进行试 验会受到很多因素限制,例如鼓筒的厚度较小,缩小 的完全几何相似模型可能无法加工。因此,设计薄壁 结构件的动力学畸变模型预测原型的固有特性具有 重要意义。

在式(12)中,当 i=0,1,…,4 时,有很多待选的畸 变相似关系

$$\lambda_{\omega} = \sqrt{\frac{\lambda_{L_{j-k_{i}}}}{\lambda_{L_{j}}\lambda_{\alpha}}\lambda_{\beta}^{k-i}}$$
(15)

根据式(7),可得 $\lambda_{l_{j-k_i}} = \lambda_E^m \lambda_{\alpha}^{o'} \lambda_{\beta}^{s'} \lambda_h^g, \lambda_{l_j} = \lambda_{\rho}^{-n} \lambda_h^{g'}$ 。因此, 薄壁结构件的畸变相似关系可以表示为

$$\lambda_{\omega} = \sqrt{\lambda_{\rm E}^{\rm m} \lambda_{\rho}^{\rm n} \lambda_{\alpha}^{\rm o} \lambda_{\beta}^{\rm s} \lambda_{\rm h}^{\rm q}} \tag{16}$$

通常在畸变相似关系中,相似因子 $\lambda_{\rm E}$ 和 λ_{ρ} 的指数 m 和 n 可以通过本构方程推导得到。然而,指数 o、 s 和 q 通常是未知的。为了确定薄壁构件精确的畸变 相似关系,提出并在理论上证明了结构参数对固有频 率敏感性值与相似因子指数的比值关系。

2.3 畸变设计准则

采用敏感性分析法确定薄壁构件精确的畸变相 似关系,所谓敏感性是指对于结构振动系统,结构特 征参数(特征值λ)对结构参数 p(质量、刚度、阻尼、 结构参数)的改变率^[16-17]。

首先给出基于敏感性分析确定畸变相似关系的 设计准则:畸变相似关系中相似因子指数的比值 k_1 : k_2 : …: k_n 可近似为结构参数对固有频率敏感性 Φ_1 : Φ_2 :…: Φ_n 的比值。即

$$\mathbf{k}_1:\mathbf{k}_2:\cdots:\mathbf{k}_n \approx \boldsymbol{\Phi}_1:\boldsymbol{\Phi}_2:\cdots:\boldsymbol{\Phi}_n \qquad (17)$$

式中: $k_1:k_2:\cdots:k_n$ 分别为相似因子 $\lambda_1,\lambda_2,\cdots,\lambda_n$ 的指数; $\Phi_1:\Phi_2:\cdots:\Phi_n$ 分别为各结构参数的敏感性值。

下面证明式(17)成立。

结构参数 α 和 β 在极限范围内 ($\lambda_{\alpha}, \lambda_{\beta} \in [1-\varepsilon, 1+\varepsilon], \varepsilon$ 为无穷小)变化的畸变模型相似因子分别为

$$\lambda_{\alpha\omega} = \frac{\omega_{\rm p}}{\omega_{\alpha t}} = \lambda_{\alpha t}^{0/2}, \lambda_{\beta\omega} = \frac{\omega_{\rm p}}{\omega_{\beta t}} = \lambda_{\beta t}^{s/2}$$
(18)

结构参数 α 和 β 的敏感性分别为

$$\Phi_{\alpha} = \frac{d\omega}{d\lambda_{\alpha}} = \frac{\Delta\omega}{\Delta\lambda_{\alpha}} = \frac{\omega_{p} - \omega_{\alpha t}}{1 - \lambda_{\alpha t}}$$

$$\Phi_{\beta} = \frac{d\omega}{d\lambda_{\beta}} = \frac{\Delta\omega}{\Delta\lambda_{\beta}} = \frac{\omega_{p} - \omega_{\beta t}}{1 - \lambda_{\beta t}}$$
(19)

因此,可得

$$\frac{\Phi_{\alpha}}{\Phi_{\beta}} = \frac{\omega_{p} - \omega_{\alpha t}}{1 - \lambda_{\alpha t}} \frac{1 - \lambda_{\beta t}}{\omega_{p} - \omega_{\beta t}} = \omega_{\alpha t} \frac{(\omega_{p} / \omega_{\alpha t} - 1)}{1 - \lambda_{\alpha t}} \frac{1 - \lambda_{\beta t}}{(\omega_{p} / \omega_{\beta t} - 1)} \frac{1}{\omega_{\beta t}} (20)$$

将式(18)代人式(20)得到

$$\frac{\Phi_{\alpha}}{\Phi_{\beta}} = \frac{\omega_{\alpha t}}{\omega_{\beta t}} \frac{\lambda_{\alpha t}^{0/2} - 1}{1 - \lambda_{\alpha t}} \frac{1 - \lambda_{\beta t}}{\lambda_{\beta t}^{s/2} - 1}$$
(21)

当相似因子在极限范围内变化时有: $\lambda_{\alpha t} \rightarrow 1, \lambda_{\beta t} \rightarrow 1$ 。此时

$$\lim_{\substack{\lambda_{\alpha} \to 1 \\ \lambda_{\beta} \to 1}} \frac{\omega_{\alpha t}}{\omega_{\beta t}} = 1, \lim_{\lambda_{\alpha} \to 1} \frac{\lambda_{\alpha t}^{s_{\alpha}} - 1}{1 - \lambda_{\alpha t}} = -\frac{0}{2}, \lim_{\lambda_{\beta} \to 1} \frac{1 - \lambda_{\beta t}}{\lambda_{\beta t}} = \frac{2}{s} \quad (22)$$

式(22)可写为

$$\frac{\Phi_{\alpha}}{\Phi_{\beta}} = \frac{\omega_{\alpha t}}{\omega_{\beta t}} \left(-\frac{0}{2}\right) \left(-\frac{2}{s}\right) = \frac{0}{s}$$
(23)

当相似比在小范围内变化时,可得

$$\frac{\Phi_{\alpha}}{\Phi_{\beta}} \approx \frac{0}{s}$$
 (24)

当薄壁结构件的固有特性受多个结构参数影响时,与上述过程同理递推即可得到式(17)的结果。因此,式(17)得证。

因此,基于上述设计准则可知薄壁结构件结构参 数对固有频率的敏感性与相似因子指数比值关系, 即:通过分析结构参数对固有频率的敏感性,即可确 定畸变相似关系中相似因子未知的指数,从而得到薄 壁构件畸变相似模型的设计准则。

最后给出适用于薄壁构件精确畸变相似关系的 确定步骤:

(1)针对典型薄壁构件,如矩形薄板等,基于本构 方程推导完全几何相似关系。

(2)假设薄壁构件的畸变相似关系。

(3)通过结构参数对固有频率的敏感性分析,根据薄壁结构件畸变相似关系的设计准则,最终确定薄壁构件畸变相似关系。

3 结论

(1)通过对薄壳微元的受力分析,推导得到适用 于薄壁构件统一形式的微分方程式(8)。

(2)根据薄壁构件统一的微分方程,建立了完全 几何相似试验模型准确预测原型动力学特性的动力 学相似关系式(14)。

(3)提出并在理论上证明了基于薄壁类构件结构 参数对固有频率的敏感性与相似因子指数比值关系 式(17),从而得到薄壁构件畸变相似试验模型的设计 准则。

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